

A Gibbs Field Model for the Microscopic Yielding of a Viscoplastic Fluid

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November 12, 2013

Outline

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- Deterministic Macroscopic Models of Viscoplastic Fluids

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- Stochastic Microscopic Model

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- Simulation Results

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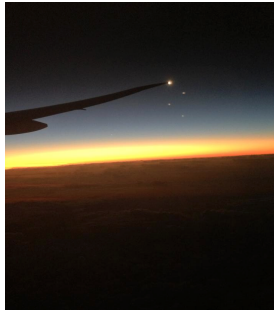
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- Summary

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- Acknowledgments

**Brief outline of the collaboration
on the transfer phenomena in viscoplastic fluids:**

BEYOND THE DATE LINE



PEOPLE INVOLVED IN THE COLLABORATION

LTN Nantes
(Experiments)



C. Castelain
(our boss, yes!)



Teo Burghilea

UC, New Zealand
(theory, modelling)

Miguel Moyers-Gonzalez

Phil Wilson

Raazesh Sainudiin

The recently founded Complex Fluids Team
(since October 2010)

Projects (past and ongoing)

(1) Hydrodynamic stability of an elasto-visco plastic material

J. Non-Newtonian Fluid Mech. 166 (2011) 515–531



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Linear stability analysis for plane-Poiseuille flow of an elastoviscoplastic fluid with internal microstructure for large Reynolds numbers

Miguel Moyers-Gonzalez^{a,*}, Teodor I. Burgehelea^b, Julian Mak^c

^a Department of Mathematics and Statistics, University of Canterbury, Private Bag 4800, Christchurch 8140, New Zealand

^b Université de Nantes, Nantes Atlantique Universités, CNRS, Laboratoire de Thermocinétique de Nantes, UMR 6607, La Chantrerie, Rue Christian Pauc, B.P. 50609, F-44306 Nantes Cedex 3, France

^c School of Mathematics, University of Leeds, Leeds LS2 9JT, UK

(2) Thermorheological properties of a Carbopol gel under shear

Journal of Non-Newtonian Fluid Mechanics 183–184 (2012) 14–24



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Thermorheological properties of a Carbopol gel under shear

Eva Weber^a, Miguel Moyers-González^b, Teodor I. Burghelca^{c,*}

^a Institute of Polymer Materials, Friedrich-Alexander-University Erlangen-Nürnberg, D-91058 Erlangen, Germany

^b Department of Mathematics and Statistics, University of Canterbury, Private Bag 4800, Christchurch 8041, New Zealand

^c Université de Nantes, Nantes Atlantique Universités, CNRS, Laboratoire de Thermocinétique de Nantes, UMR 6607, La Chantrerie, Rue Christian Pauc, B.P. 50609, F-44306 Nantes Cedex 3, France

(3) Unsteady flows of a Carbopol gel (being reviewed at JNNFM)

Unsteady laminar flows of a Carbopol[®] gel in the presence of wall slip

Antoine Poumaere

*LUNAM Université, Université de Nantes, CNRS, Laboratoire de Thermocinétique, UMR 6607, La Chantrerie, Rue Christian Pauc, B.P. 50609,
F-44306 Nantes Cedex 3, France*

Miguel Moyers-González

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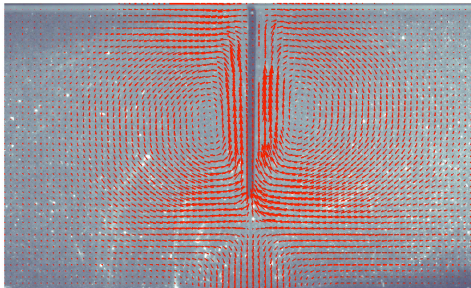
Cathy.Castelain@univ-nantes.fr

Teodor Burghilea

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F-44306 Nantes Cedex 3, France*

(4) The Landau-Levich coating problem with a yield stress fluid (in preparation)

(Miguel Moyers-Gonzalez, Phil Wilson, Cathy Castelain, Teo Burghlea)

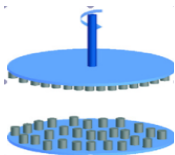


The QUEST:
**Understand the flows of viscoplastic materials
through a multi scale approach**

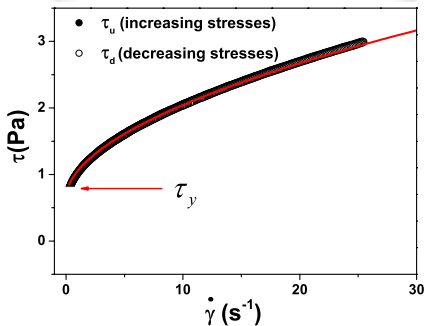
Viscoplastic Material:

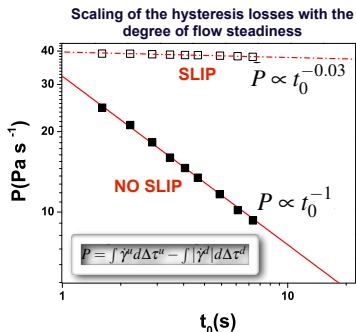
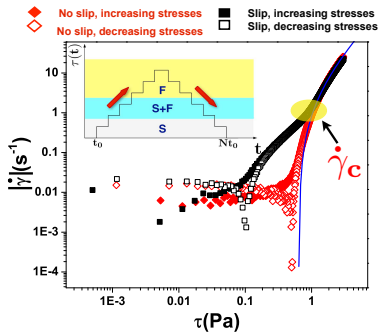
**A material that behaves as a solid at low applied stresses and as a fluid
beyond a threshold value of the applied stress**

Yielding and flow of a gel at a macroscopic scale (Teo, Cathy, Miguel, Phil)



The common “religion” of yielding: H-B





An unsteady rheological test takes us from paradise to hell in several ways...

- Rheological hysteresis

- Gradual Solid-Fluid transition coupled to the wall slip

- NO STEADY STATE REACHABLE IN THE PRESENCE OF WALL SLIP

The **rheological hysteresis** for a Carbopol gel is a new (and still subject of some controversy) observation:

Rheol Acta
DOI 10.1007/s00397-009-0365-9

ORIGINAL CONTRIBUTION

The solid–fluid transition in a yield stress shear thinning physical gel

Andreas M. V. Putz · Teodor I. Burghelca

Received: 28 May 2008 / Accepted: 23 April 2009
© Springer-Verlag 2009

Other groups observed it as well

Soft Matter

Dynamic Article Links 

Cite this: DOI: 10.1039/c1sm05607g

www.rsc.org/softmatter

PAPER

From stress-induced fluidization processes to Herschel-Bulkley behaviour in simple yield stress fluids

Thibaut Divoux,[†] Catherine Barentin[†] and Sébastien Manneville[†]

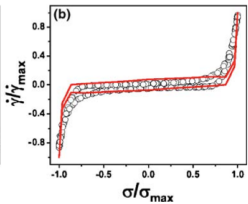
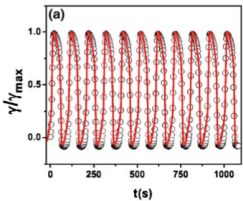
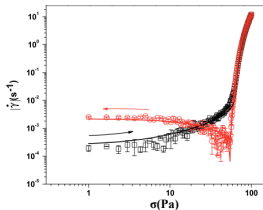
Received 6th April 2011, Accepted 25th May 2011
DOI: 10.1039/c1sm05607g

RECONCILE THE "HEAVEN" AND THE "HELL": THE POOR MAN APPROACH AND BEYOND

$$\frac{d\Phi}{dt} = R_d(\Phi, t, \Gamma) + R_r(\Phi, t, \Gamma) + \delta$$

$$R_d(\Phi, t, \Gamma) = -K_1\Gamma\Phi, \quad \Gamma = \frac{\sigma}{\sigma_2}$$

$$R_r(\Phi, t, \Gamma) = f(\Gamma)\Phi(1 - \Phi) \quad \text{with } f(\Gamma) = K_r \left[1 - \tanh\left(\frac{\Gamma - 1}{w}\right) \right]$$



**Can this novel yielding picture be validated at a microscopic scale
(and automatically validated from a thermodynamical standpoint)?**

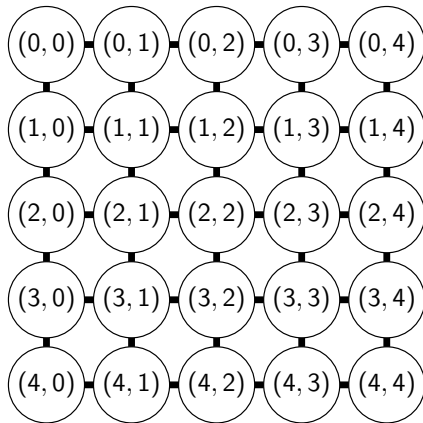
Yielding and flow of a gel at a microscopic scale:

**A Gibbs field model for the microscopic yielding
of a gel**

Raaz, we are all ears!

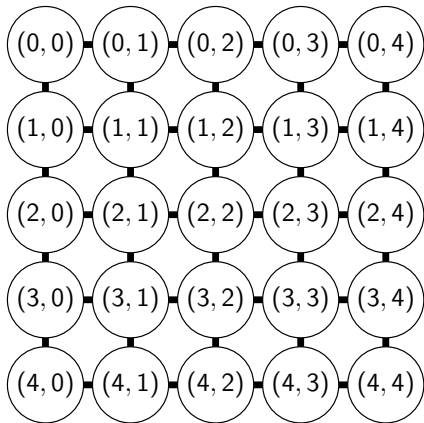
Stochastic Microscopic Model – 1

- \mathcal{G} = regular graph (2D toroidal square lattice)



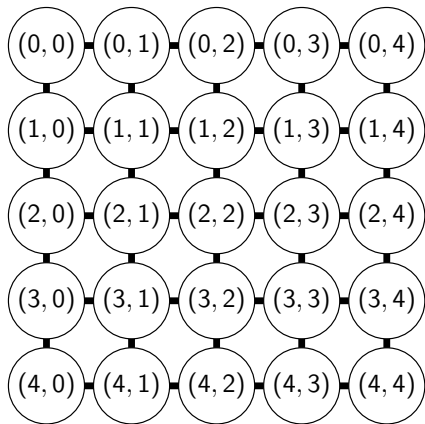
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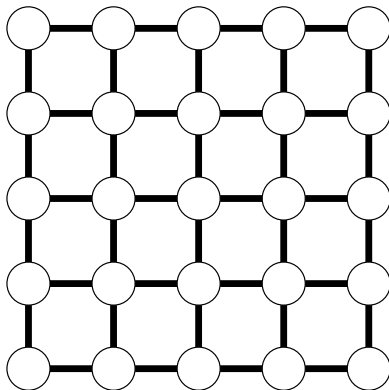
Stochastic Microscopic Model – 1

- \mathcal{G} = regular graph (2D toroidal square lattice)
- $\mathbb{S}_n = \{0, 1, 2, \dots, n - 1\}^2 =$ set of nodes or sites
- $\mathbb{E}_n \subset \mathbb{S}_n^2 =$ set of edges between pairs of sites



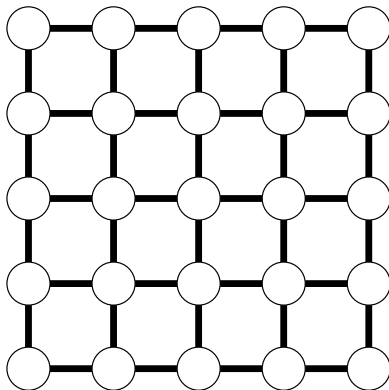
Microscopic Models (Stochastic) – 2

- each site $s \in \mathbb{S}_n$ represents a polymer molecule



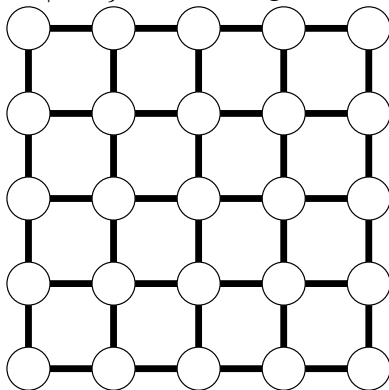
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- each edge $\langle s, t \rangle \in \mathbb{E}_n$ represents a potential bond between neighbouring molecules at sites s and t .



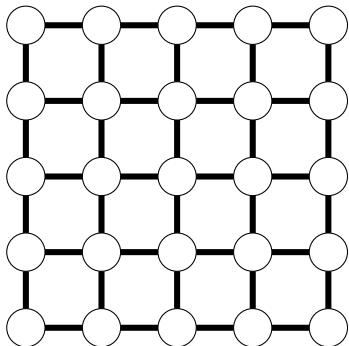
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- $N_s = \{t : |t - s| = 1\}$ denotes neighbours of a site s



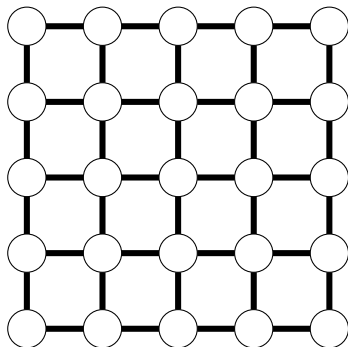
Microscopic Models (Stochastic) – 3

- $x(s) \in \Lambda = \{0, 1\}$ denote the phase of the molecule at site s



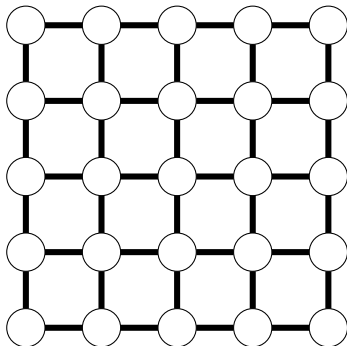
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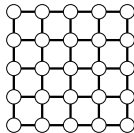
- $x(s) \in \Lambda = \{0, 1\}$ denote the phase of the molecule at site s
- phase 0 corresponds to being unbonded or isolated from all its neighbours
- phase 1 corresponds to possibly being bonded with at least one of its neighbours



Phases at Sites and Bonds

We say a bond exists between molecules at sites s and t and denote it by

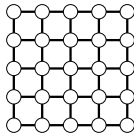
$$y(\langle s, t \rangle) = \begin{cases} 1 & \text{if } t \in N_s \text{ and } x(t)x(s) = 1 \\ 0 & \text{if } t \in N_s \text{ and } x(t)x(s) = 0 \end{cases}$$



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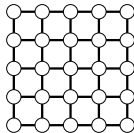


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- every site configuration $x \in \mathbb{X}_n := \Lambda^{\mathbb{S}^n}$ has an associated bond configuration $y \in \mathbb{Y}_n := \Lambda^{\mathbb{E}^n}$
- Note that $Y(x) : \mathbb{X}_n \rightarrow \mathbb{Y}_n$ is neither injective nor surjective

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where

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$$V_{\langle s, t \rangle}(x) = -\frac{B}{k} x(s)x(t) = \begin{cases} 0 & \text{if } (x(s), x(t)) = (0, 0) \\ 0 & \text{if } (x(s), x(t)) = (1, 0) \\ 0 & \text{if } (x(s), x(t)) = (0, 1) \\ -\frac{B}{k} & \text{if } (x(s), x(t)) = (1, 1) \end{cases}$$

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where $\langle s, t \rangle$ is the two element clique with $t \in N_s$ and B is the internal energy of a bond-pair (between two polymer molecules).

Total, Internal and Free Energy – 1

Thus, the energy function corresponding to this potential is

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$$\begin{aligned} \mathcal{E}(x) &= \sum_C V_C(x) = \sum_{s \in \mathbb{S}_n} V_{\{s\}}(x) + \sum_{\langle s, t \rangle \in \mathbb{E}_n} V_{\langle s, t \rangle}(x) \\ &= \frac{1}{k} \left(-B \sum_{\langle s, t \rangle \in \mathbb{E}_n} x(s)x(t) + (\sigma - \sigma_0) \sum_{s \in \mathbb{S}_n} x(s) \right) \end{aligned}$$

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Let the expectation of a function $g : \mathbb{X}_n \rightarrow \mathbb{R}$, w.r.t. π , be

$$\langle g \rangle := \sum_{x \in \mathbb{X}_n} g(x) \pi(x)$$

Total, Internal and Free Energy – 2

then the *internal energy* of the system is

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$$\mathcal{U} = -T^2 \frac{\partial}{\partial T} \left(\frac{\mathcal{F}}{T} \right)$$

Energy \mathcal{E} and Gibbs Distribution π over Site Configurations

Let $\mathcal{E}(x)$, the energy of a site configuration x be

$$\mathcal{E}(x) = \frac{1}{k} \left(-B \sum_{\langle s,t \rangle \in \mathbb{E}_n} x(s)x(t) + (\sigma - \sigma_0) \sum_{s \in \mathbb{S}_n} x(s) \right)$$

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then the probability distribution of interest on the site configuration space \mathbb{X}_n is

$$\pi(x) = \frac{1}{Z_T} \exp \left(-\frac{1}{T} \mathcal{E} \right),$$

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where Z_T is the normalising constant or partition function

$$Z_T = \sum_{x \in \mathbb{X}_n} \exp \left(-\frac{1}{T} \mathcal{E} \right).$$

Two Informative Statistics of a Configuration $x - 1$

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The number of “bondable” polymers

$$a(x) = \sum_{s \in \mathbb{S}_n} x(s),$$

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$$a(x) = \sum_{s \in \mathbb{S}_n} x(s),$$

The average number of bondable polymers per site is

$$\bar{a} := |\mathbb{S}_n|^{-1} \langle a \rangle = n^{-2} \sum_{x \in \mathbb{X}_n} a(x) \pi(x) = n^{-2} \sum_{x \in \mathbb{X}_n} \left(\sum_{s \in \mathbb{S}_n} x(s) \right) \pi(x)$$

Two Informative Statistics of a Configuration x – 2

The number of bonds

$$b(x) = \sum_{\langle s,t \rangle \in \mathbb{E}_n} y(\langle s,t \rangle) = \sum_{\langle s,t \rangle \in \mathbb{E}_n} x(t)x(s)$$

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and the average number of bonds per edge (or pair of sites) is

$$\begin{aligned} \bar{b} &:= |\mathbb{E}_n|^{-1} \langle b \rangle = (2n^2)^{-1} \sum_{x \in \mathbb{X}_n} b(x) \pi(x) \\ &= \frac{n^{-2}}{2} \sum_{x \in \mathbb{X}_n} \left(\sum_{\langle s,t \rangle \in \mathbb{E}_n} x(t)x(s) \right) \pi(x) \end{aligned}$$

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$$\begin{aligned} \bar{b} &:= |\mathbb{E}_n|^{-1} \langle b \rangle = (2n^2)^{-1} \sum_{x \in \mathbb{X}_n} b(x) \pi(x) \\ &= \frac{n^{-2}}{2} \sum_{x \in \mathbb{X}_n} \left(\sum_{\langle s, t \rangle \in \mathbb{E}_n} x(t)x(s) \right) \pi(x) \end{aligned}$$

One of our primary interests is to study \bar{a} and \bar{b} as a function of externally applied stress σ .

Gibbs Sampler – 1

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We want to simulate configurations $x \in \mathbb{X}_n$ that are distributed according to π (for fixed $\sigma, \sigma_0, k, B, T$)

- We construct an \mathbb{X}_n -valued Markov chain $\{X_m\}_{m \geq 0}$, where $(X_m(s), s \in \mathbb{S}_n)$ and $X_m(s) \in \Lambda$, with stationary distribution π , i.e.,

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Next we derive the algorithm for the Gibbs sampler.

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where

$$\theta = \exp\left(-\frac{1}{T} \frac{1}{k} \left(B \sum_{t \in N_s} x(t) - (\sigma - \sigma_0) \right)\right)$$

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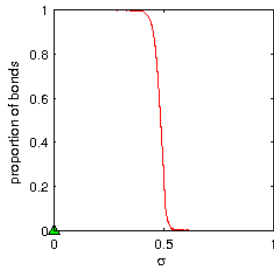
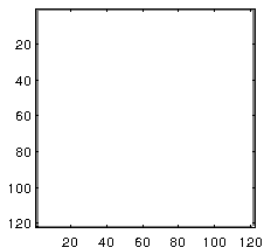
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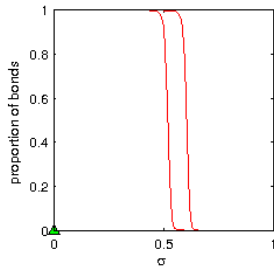
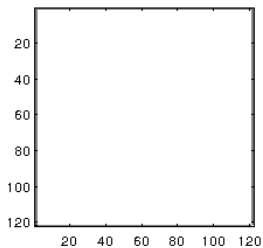
When $B = 0$

$K = 0.015625$ $B = 0$ Hits/Site = 10 time = 144000



$$B \neq 0$$

$K = 0.015625$ $B = 0.015625$ Hits/Site = 1 time = 14400



The End

- Many thanks to:
 - Teo and Cathy for inviting me here
 - UC for research travel funds
 - Thanks for your Attention :)